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direction of the light. Because of the circularly symmetric nature of the diverging light beam, the transverse polarization generates a differential absorption. For example, if the atomic polarization has a component along the x direction, light that has a component in the x direction will see reduced absorption, while light that has a component in the –x direction will see increased absorption. As the atomic polarization vector precesses, therefore the spatial profile of the diverging light field will be modified accordingly.

FIG. 2 illustrates a schematic of the operation of a diverging beam magnetometer 100 in accordance with one embodiment of the present invention. The instrument is placed in the vicinity of the DC magnetic field to be sensed,  $B_0$ , 110. Atoms, 120 are polarized along the average direction of propagation of the incident light field, 130. RF coils, 140, 15 generate an oscillating magnetic field at the Larmor frequency of atoms in a direction perpendicular to the magnetic field. This causes the atomic polarization to precess at the drive frequency about the magnetic field B<sub>o</sub>, as shown by the precessing vector, 150. The transverse component of the 20 atomic polarization, 160, is detected by monitoring the absorption of the edges of the diverging light beam, 170. By subtracting the signals coming from opposite sides of the light beam, the magnitude of the precessing transverse atomic polarization can be determined.

FIG. 3 illustrates a schematic of the operation of a diverging beam gyroscope 200 in accordance with one embodiment of the present invention. The instrument is placed in a wellcontrolled, known, uniform magnetic field, Bo, 205. The spin of alkali atoms (such as <sup>133</sup>Cs, <sup>87</sup>Rb, or <sup>39</sup>K), **210** is polarized 30 via optical pumping with a component along the average direction of propagation of the incident light field, 215. The spin of the nuclei of noble gas atoms, 220 (129Xe, 3He or equivalent) is polarized with a component along the average direction of propagation of the incident light field via spin 35 exchange collisions with the alkali atoms. One set of RF coils, 225, generate an oscillating magnetic field,  $\square B_1$ , 230, in a direction perpendicular to the magnetic field, B<sub>o</sub>, 205 at the Larmor frequency of noble gas atoms 220. This causes the nuclear polarization to precess at the drive frequency of the 40 RF coils 225 about the magnetic field B<sub>0</sub> 205. A second AC magnetic field  $\square B_2$  245 is applied parallel to the static field  $B_a$ 205, at the Larmor frequency of the alkali atoms 210. This field causes the alkali atom spins to precess, 250, about the total field created by the static field, B<sub>o</sub> 205, and the precess-45 ing polarization of the noble gas nuclei, 220. The transverse component of the atomic polarization is detected by monitoring the absorption of the edges of the diverging light beam, 255. By subtracting the signals, 260, coming from opposite sides of the light beam, the magnitude of the precessing 50 transverse atomic polarization, and in turn the Larmor frequency of both the alkali and the noble gas species. The angular rate of the vessel, can be determined from the induced shift in the measured noble gas Larmor frequency.

FIG. 4 illustrates a schematic of an implementation of the diverging beam magnetometer or gyroscope 300 in accordance with one embodiment of the present invention. In this implementation, the diverging components of the laser beam 310 are reflected by angled walls inside the alkali vapor cell 330 to form counter-propagation probe beams 340. These 60 beams are reflected a second time off the opposite wall and the power is detected by photodiodes 350 placed on the base plate 360 with the laser 370.

FIG. 5 illustrates a schematic of the basic operation of a co-magnetometer nuclear magnetic resonance ("NMR") gyroscope as a co-magnetometer. The diverging light beam 410 polarized the alkali species 420 in the nominal direction

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of the light propagation. The noble gas species 430 is subsequently polarized via spin-exchange collisions with the alkali atoms. A longitudinal magnetic field B<sub>0</sub> 440 is applied with a magnitude that largely cancels the field due to the noble gas as seen by the alkali atoms. Under rotation, the noble gas spins 430 become misaligned with the longitudinal field 440 and the spin orientation rotates slightly about the component of the longitudinal field 440 perpendicular to the noble gas spin orientation. The rotation of the noble gas spin 430 causes a small transverse field  $B_{tran}$  450 seen by the alkali atoms 420. This transverse component 550 causes the orientation of the alkali species, 460, to change slightly. The orientation change of the alkali species causes a differential absorption of the counter-propagating light fields 470. The difference in absorption is measured by two photodetectors 480. The difference in the signals measured by the photodetectors 480 is proportional to the instrument rotation rate.

FIG. 6 illustrates a schematic of a cross-sectional view of a compact nuclear magnetic gyroscope 500 ("NMRG") in accordance with one embodiment of the present invention. A multilayer magnetic shield 510 (here 3 layers) is used to suppress the external magnetic field 520 (e.g., Earth's field, fields created by adjacent electrical components, and other environmental fields) by over six orders of magnitude. Inside the shields, a set of 3-axis coils 530 are used to create a very precise static magnetic field B<sub>o</sub>. Additionally, the coils 530 are also used to compensate for residual magnetic fields (external or internal to the shields) that may exist in the area of the NMR cell **540**. The light from the laser on the base-plate is circularly polarized by a wave plate and transmitted through the NMR cell **540**. The light is then reflected off the angled cell walls back onto photodiodes on the base plate 550. Two small flex circuits 570 provide paths for electrical signals for both the base plate 550 and the 3-axis coils 530 to flow between the interior of the shield 510 and the exterior.

FIG. 7(a) and FIG. 7(b) illustrate a schematic of the exterior 600 and cross-sectional view of magnetic shields 610 in accordance with one embodiment of the present invention. While a 3-layer shield is shown for illustrative purposes, the number of layers will vary depending on the desired shielding factor and the particular application. The shields may be machined and welded from a high-permeability material such as mumetal. The nuclear magnetic resonance gyroscope is to be positioned at the center of the shields and the flex circuits from the coils and the base plate is to be threaded through the holes 620 of the shield caps 630. The spacers 640 that are used to separate the shield layers are machined from a nonmagnetic material.

FIG. 8(a) and FIG. 8(b) illustrate a schematic of a twolayer flex circuit fabricated on a planar substrate and a flex circuit wrapped around a cylindrical holder to form a set of three-axis coils 700 in accordance with one embodiment of the present invention. A set of three-axis coils 700 are used to generate the static and oscillating magnetic fields (B<sub>o</sub> **205**  $\square B_1$  **230** and  $\square B_2$  **245** in FIG. **3**) as well as for compensation of residual fields. The coils 710 are fabricated as highly conductive traces (e.g. metal) on a two-layer flexible planar substrate 720, such as polyimide. The planar substrate is then wrapped around a machined cylindrical holder 730 and the conductive pads on two opposite sides of the substrate are brought in contact and affixed to each other (e.g. soldered) to form a set of three-axis coils 700. The tail end of the flex of the coils 740 contains bonding pads, which allows the coils 700 to be connected to external circuitry once the flex has been threaded through the shields. Once assembled, both the length and the diameter of the coil structure are on the order of 2 to 10 millimeters. The coil for the longitudinal magnetic